

## Effect of Axial Ratio Changes on the Elastic Moduli and Grüneisen $\gamma$ for Lower Symmetry Crystals

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Gerlich has shown that Sheard's model for calculating mode  $\gamma$ 's from hydrostatic pressure derivatives of the elastic moduli of hcp Mg and Cd yields Grüneisen  $\gamma$ 's at both high and low temperatures that are in good agreement with the  $\gamma$ 's derived from thermal-expansion measurements. For hcp Ti and Zr, however, large differences arise, primarily from very small values for  $dC_{44}/dP$ . It is proposed that these small values are caused by the changes in  $c/a$  ratio with hydrostatic pressure because of a large dependence of  $C_{41}$  on the  $c/a$  ratio. The disagreement with thermal-expansion data can be removed by taking into account the difference in  $d(c/a)/dV$  between hydrostatic-pressure and thermal-expansion conditions. The effect of  $\Delta(c/a)$  is not found in tetragonal  $\text{TiO}_2$ , rutile, where  $\bar{\gamma}_H$  is in excellent agreement with the thermal expansion  $\gamma_\infty$ .

### INTRODUCTION

The model of Sheard<sup>1,2</sup> for calculating the Grüneisen  $\gamma$  from measured values of the hydrostatic-pressure derivatives of the elastic moduli of cubic crystals has been extended by Gerlich<sup>3</sup> to hcp and tetragonal crystals. This extension does not consider the effects of the changes in  $c/a$  ratio with pressure on the mode frequencies. In this paper we present evidence that the effects of changing the axial ratios, expressed as  $\Delta(c/a)$ , are quite large in certain hcp metals and contribute to negative mode  $\gamma$ 's and large deviations between average gammas computed from thermal expansion and those derived from measured values of  $dC_{ij}/dP$ .

### STATEMENT OF THE PROBLEM

The equations for analyzing the  $\Delta(c/a)$  effect are stated as follows:

$$\gamma^p(q) = [\gamma^p(q)]_{c/a} - \left( \frac{\partial \ln \omega_p(q)}{\partial \ln(c/a)} \right)_V \frac{d \ln(c/a)}{d \ln V}, \quad (1)$$

where  $\gamma^p(q)$  is the mode gamma for hydrostatic pressure as derived from the  $dC_{ij}/dP$  using Gerlich's<sup>3</sup> Eq. (11). The first term on the right is the pure volume change contribution, and the second term is the  $\Delta(c/a)$  contribution. A problem arises because  $d \ln(c/a)/d \ln V$  differs under thermal-expansion and hydrostatic-pressure conditions, respectively, as follows:

$$\frac{d \ln(c/a)}{d \ln V} = \left( \frac{\partial \ln(c/a)}{\partial \ln V} \right)_P = \frac{\alpha_{||} - \alpha_{\perp}}{\alpha_V} \quad (2)$$

$$\frac{d \ln(c/a)}{d \ln V} = \left( \frac{\partial \ln(c/a)}{\partial \ln V} \right)_T = \frac{\beta_{||} - \beta_{\perp}}{\beta_V}, \quad (3)$$

where  $\alpha$  and  $\beta$  are used for thermal-expansion coefficient and isothermal compressibility, respectively. The subscripts refer to the linear parameter values parallel and perpendicular, respectively, to the  $c$  axes of either hexagonal or tetragonal crystals, and  $\alpha_V$  and  $\beta_V$  are the volume parameters. If

$$\left( \frac{\partial \ln \omega_p(q)}{\partial \ln(c/a)} \right)_V$$



TABLE I. Hydrostatic pressure derivatives of elastic moduli and the anisotropy parameters for Ti, Zr, Mg, and Cd at 300°K.

	$dC_{11}/dP$	$dC_{33}/dP$	$dC_{44}/dP$	$dC_{66}/dP$	$(\beta_{11}-\beta_{\perp})/\beta_V$	$(\alpha_{11}-\alpha_{\perp})/\alpha_V^a$
Ti	5.01	4.88	0.52	0.45	0.013	-0.144(a), 0.045(b)
Zr	3.93	5.49	-0.22	0.26	-0.049	0.136(a), 0.059(b)
Mg	6.11	7.22	1.58	1.36	0.013	0.019
Cd	9.29	7.26	2.38	2.59	0.660	0.361

<sup>a</sup> Ti (a) Ref. 12, Ti (b) Ref. 13, Zr (a) Ref. 14, Zr (b) Ref. 15, Mg and Cd Ref. 10.

is a significant quantity and  $(\alpha_{11}-\alpha_{\perp})/\alpha_V$  differs from  $(\beta_{11}-\beta_{\perp})/\beta_V$ , the average mode  $\gamma$ 's for low and high temperatures ( $\tilde{\gamma}_L$  and  $\tilde{\gamma}_H$ ), as defined by Gerlich, will necessarily differ from  $\gamma_L(\alpha_V)$  and  $\gamma_H(\alpha_V)$  computed from thermal-expansion data.

### EVIDENCE FOR THE $\Delta(c/a)$ EFFECT

Experimental evidence that the  $\Delta(c/a)$  effect must be considered arises when a comparison is made of the  $dC_{ij}/dP$  values for hcp Ti<sup>4</sup> and Zr.<sup>5</sup> These two metals are exceptionally similar in many physical and mechanical properties. In regard to elastic properties,<sup>6</sup> they differ primarily in the values of the  $C_{44}$  shear modulus and in the linear compressibility perpendicular to the hexagonal axes,  $\beta_{\perp}$ . As a consequence of the latter difference,

$$d(c/a)/dP = (c/a)(\beta_{\perp}-\beta_{11}) \quad (4)$$

is negative for Ti and is positive for Zr.

The hydrostatic-pressure derivatives of the single-crystal elastic moduli of Ti, Zr, Mg, and Cd<sup>7,8</sup> are listed in Table I. All the data are taken from ultrasonic velocity measurements at temperatures near 25°C and represent adiabatic pressure derivatives. From these data, the pressure derivatives of the shear stiffnesses  $dC_{44}/dP$  and  $dC_{66}/dP$  appear to decrease with the  $c/a$  ratio of the crystal at 1-bar pressure. The major difference between Ti and Zr appears in  $dC_{44}/dP$ , where we find a negative value for Zr.

To relate the measured  $dC_{ij}/dP$  to the volume and  $c/a$  changes, separately, we use the following equations

$$\begin{aligned} \frac{dC_{ij}}{dP} &= \left( \frac{\partial C_{ij}}{\partial P} \right)_{c/a} - \left( \frac{\partial C_{ij}}{\partial(c/a)} \right)_V \left( \frac{d(c/a)}{dP} \right) \\ &= -\beta_V C_{ij} \left( \frac{\partial \ln C_{ij}}{\partial \ln V} \right)_{c/a} + (c/a)(\beta_{\perp}-\beta_{11}) \left( \frac{\partial C_{ij}}{\partial(c/a)} \right). \end{aligned} \quad (5)$$

For cubic metals the measured values of  $dC_{ij}/dV$  are negative in all cases, ( $dC_{ij}/dP > 0$ ), and we can reasonably presume that  $(\partial C_{ij}/\partial V)_{c/a}$  for Ti and Zr are also negative. We then conclude that the negative value for  $dC_{44}/dP$  in Zr, where  $\beta_{\perp} > \beta_{11}$ , arises from a negative value for  $[\partial C_{44}/\partial(c/a)]_V$ . To estimate the relative values for volume and  $c/a$  contributions to  $dC_{ij}/dP$ , it appears

reasonable to compute  $(\partial C_{ij}/\partial V)_{c/a}$  and  $[\partial C_{ij}/\partial(c/a)]_V$  by assuming that these two unknown factors are the same in Zr and Ti. The computed values of the partial derivatives and the volume and  $c/a$  contributions to each  $dC_{ij}/dP$  are listed in Table II. The conclusions from this approach are that the change in  $c/a$  with pressure has a larger effect on the  $C_{44}$  of Zr than does the volume change, and also contributes significantly to  $dC_{44}/dP$  and  $dC_{66}/dP$  in both metals.

Two external factors indicate that the quantities derived from the above procedure are realistic. One is that the value for  $dC_{44}/d(c/a)$  is very near the value derived by Cousins from calculations of the electrostatic contribution to  $C_{44}$  of hcp metals.<sup>9</sup> These calculations give

$$\left( \frac{dC_{44}^E}{d(c/a)} \right)_V = -\frac{Z^2}{a_0^4} (26.4 \times 10^{12} \text{ dyn/cm}^2), \quad (7)$$

where  $Z$  is the effective valence, and  $a_0$  is the ion separation in the basal plane in Å units. Assuming  $Z=4$  for Zr and Ti, Eq. (7) gives  $[\partial C_{44}^E/\partial(c/a)]_V = -3.88$  and  $-5.57 \times 10^{12}$  dyn/cm<sup>2</sup> for Zr and Ti, respectively, whereas our common value is  $-6.5 \times 10^{12}$  dyn/cm<sup>2</sup>. This near agreement suggests that our assumptions in deriving  $dC_{44}/d(c/a)$  are reasonably valid, and that the large effect of  $\Delta(c/a)$  on  $C_{44}$  is caused primarily by the change in electrostatic energy contribution.

The other factor that lends confidence to the procedure is that the difference between the values of  $C_{44}$  in Ti and in Zr is quite large<sup>6</sup> and can be reasonably accounted for from the  $(\partial C_{44}/\partial V)_{c/a}$  and  $[\partial C_{44}/\partial(c/a)]_V$  contributions. The total observed difference is  $0.145 \times 10^{12}$  dyn/cm<sup>2</sup>. The volume difference can account for approximately  $0.110 \times 10^{12}$ , and the  $c/a$  difference accounts for  $0.023 \times 10^{12}$  dyn/cm<sup>2</sup>, when the derived values are used for the partial derivatives.

### Grüneisen $\gamma$ Calculations in Ti and Zr

Some results of the  $\tilde{\gamma}_L$  and  $\tilde{\gamma}_H$  calculations, using Eq. (2), Gerlich's computer program,<sup>3</sup> and measured  $dC_{ij}/dP$ , are listed in Table III. The results for Mg and Cd were obtained and reported by Gerlich.<sup>3</sup> The excellent agreement between  $\tilde{\gamma}_L$  and  $\tilde{\gamma}_L(\alpha_V)$  and the good agreement between  $\tilde{\gamma}_H$  and  $\tilde{\gamma}_H(\alpha_V)$  for Mg and Cd<sup>10</sup> serve to further verify the validity of the model